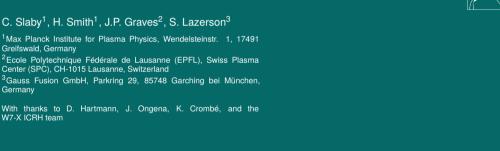
This work has been carried out within the framework of the FUROfusion Consortium funded by the European Union via the Euratom Research and Training Programme (Grant Agreement No 101052200 - EUROfusion). Views and opinions counter there of the authorial only and do not necessarily reflect these of the European Union or the mean Commission. Neither the European Union nor the European Commission can be held responsible for them

Greifswald, Germany

Germany

W7-X ICBH team



ICRH simulations for the Wendelstein 7-X stellarator

EUROfusion

Center (SPC), CH-1015 Lausanne, Switzerland







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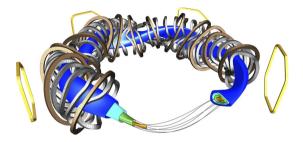
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The ICRH system at Wendelstein 7-X



- W7-X is a large, superconducting, optimized stellarator being operated at Greifswald, Germany¹
- the plasma can be heated via ECRH, NBI, and ICRH

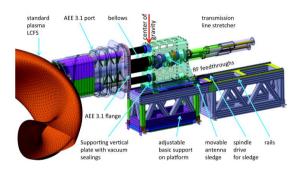


¹T. Sunn Pedersen et al., *Nature Communications* **7**, 13493 (2016)

²J. Ongena et al., *Phys. Plasmas* **21**, 061514 (2014)

The ICRH system at Wendelstein 7-X





- W7-X is a large, superconducting, optimized stellarator being operated at Greifswald, Germany¹
- the plasma can be heated via ECRH, NBI, and ICRH

The ICRH system:

- two-strap antenna whose shape is adjusted to the LCFS of the standard configuration²
- \cdot so far: 0- and $\pi\text{-phasing}$ of the antenna straps
- operational frequencies: f = 25 or 37.5 MHz
- · typical heating scenarios at 2.5 T:
 - $\,$ H-minority heating at 37.5 $\rm MHz$
 - heating of ^{3}He (minority- or 3-ion scheme) at 25 $\rm MHz$
- other uses: low-field start-up and wall conditioning
- ¹T. Sunn Pedersen et al., *Nature Communications* 7, 13493 (2016)
- ²J. Ongena et al., *Phys. Plasmas* **21**, 061514 (2014)

Models of different complexity for describing ICRH physics in W7-X



- goal: give an overview of the tools currently used for ICRH modelling at W7-X
- SCENIC³ is the workhorse for modelling ICRH in W7-X numerically
 - hot-plasma model
 - coupled to Fokker-Planck solver VENUS-LEVIS⁴ to include contributions of fast ions to dielectric tensor
 - numerically expensive \longrightarrow cheaper and simplified tools are desirable
- new overview tool has been developed⁵ for quickly plotting locations of resonances, cut-offs, and evanescence layers \longrightarrow complements the SCENIC modelling
 - cold-plasma model
 - interfaced with the VMEC webservice⁶ for the magnetic equilibrium (magnetic field and coordinate transformations)
 - fast and flexible
 - can be used for initial analysis of an experiment and can inform planning of future experiments
- ³M. Jucker et al., *Comput. Phys. Commun.* **183**, 912-925 (2011)
- ⁴D. Pfefferlé et al., Comput. Phys. Commun. 185, 3127-3140 (2014)
- ⁵C. Slaby, An overview tool for ICRH physics in Wendelstein 7-X, *IPP Report* (2024)
- ⁶M. Grahl et al., IEEE Transactions on Plasma Science 46, 1114 (2018)

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Wendelstein 7×

A simple cold-plasma model

• we use the cold-plasma dielectric tensor as derived by e.g. T.H. Stix⁷ (for $\mathbf{B} = B\hat{\mathbf{z}}$)

$$\underline{\varepsilon} = \begin{pmatrix} S & iD & 0\\ -iD & S & 0\\ 0 & 0 & P \end{pmatrix}$$
(1)

with

$$S = 1 + \sum_{\alpha} \frac{\omega_{p\alpha}^2}{\Omega_{\alpha}^2 - \omega^2} \qquad D = \sum_{\alpha} \frac{\omega_{p\alpha}^2 \Omega_{\alpha}}{\omega \left(\Omega_{\alpha}^2 - \omega^2\right)} \qquad P = 1 - \sum_{\alpha} \frac{\omega_{p\alpha}^2}{\omega^2}$$
(2)

and

$$\omega_{\mathrm{p}\alpha} = \sqrt{\frac{q_{\alpha}^2 n_{\alpha}}{m_{\alpha} \varepsilon_0}} \qquad \Omega_{\alpha} = \frac{q_{\alpha} B}{m_{\alpha}}$$
(3)

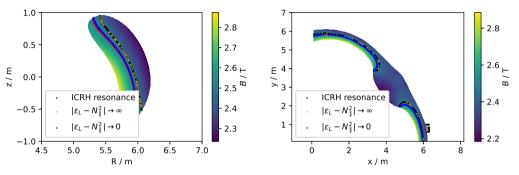
· can easily be derived from multi-fluid model, start with linearized momentum equation for species α , set $T_{\alpha} = 0$, which gives an expression

$$\mathbf{j} = \sum_{\alpha} q_{\alpha} n_{\alpha} \mathbf{u}_{\alpha} = \ldots = \underline{\sigma} \mathbf{E}$$
(4)

to be combined with Maxwell's equations

⁷T.H. Stix, Waves in Plasmas, American Institute of Physics, New York (1992)

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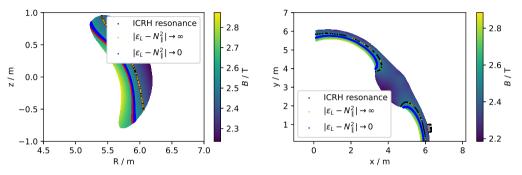
⁴He-(H) minority-heating scheme

 \cdot density fractions ($n_{\rm e,0} = 5.0 \cdot 10^{19} {
m m}^{-3}$)

species α	e^{-}	$^{4}\mathrm{He}$	Η
$n_{lpha}/n_{ m e}$	1.0	0.475	0.05

- L-cutoff and evanescence layer form mode-conversion layer on the HFS of the ICRH resonance
- locations of L-cutoff and resonance not exactly the same
- with increasing minority-density fraction: MC-layer moves further to the HFS and away from the resonance location
- $\Rightarrow\,$ heating and fast-ion generation will become less efficient





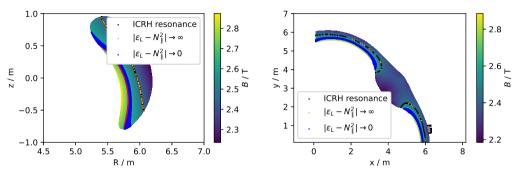
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species α	e^{-}	$^{4}\mathrm{He}$	Η
$n_lpha/n_{ m e}$	1.0	0.45	0.10

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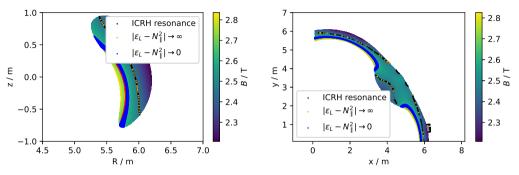
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species α	e^{-}	$^{4}\mathrm{He}$	Η
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species α	e^{-}	$^{4}\mathrm{He}$	Η
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- magnetic configuration impacts the location of the ICRH resonance
- while in the standard configuration the resonance leaves the plasma, a resonance is present at all toroidal angles in the **low-mirror configuration**
- ⇒ implications regarding power absorption, heating, and fast-ion generation





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The SCENIC code

ANIMEC

the numerical tool used to model ICRH physics in 3D equilibria is the SCENIC code⁸ developed at EPFL

fast particle effects: **ICRH** wave updated dielectric fast particle effects: tensor current / anisotropic $\left| \left\langle \Delta \mathbf{v}_{\perp}^{2} \right\rangle = \tau^{2} \gamma \frac{Z_{\alpha}^{2}}{m_{\alpha}^{2}} \left| \mathbf{E}^{+} \mathbf{e}^{-\mathrm{i}\psi} J_{n-1} \left(\frac{\mathbf{k}_{\perp} \mathbf{v}_{\perp}}{\Omega_{\alpha}} \right) + \mathbf{E}^{-} \mathbf{e}^{+\mathrm{i}\psi} J_{n+1} \left(\frac{\mathbf{k}_{\perp} \mathbf{v}_{\perp}}{\Omega_{\alpha}} \right) \right|^{2}$ pressure $\dot{\mathbf{X}} = \mathbf{v}_{\parallel} \frac{\mathbf{B}^{\star}}{\mathbf{B}^{\star}_{\parallel}} + \frac{\mathbf{E}^{\star} imes \mathbf{b}}{\mathbf{B}^{\star}_{\parallel}}$ SCENIC implements iterative procedure of three (coupled) codes

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ANIMEC: anisotropic equilibrium / otherwise similar to VMEC

 $\dot{\mathbf{v}}_{\parallel} = rac{q}{m} rac{\mathbf{B}^{\star} \cdot \mathbf{E}^{\star}}{\mathbf{B}^{\star}}$

 $\nabla^2 \mathbf{A} + k_0^2 \hat{\epsilon} \cdot \mathbf{A} + \mathrm{i} k_0 \hat{\epsilon} \cdot \nabla \phi = -\frac{4\pi}{2} \mathbf{j}_{\mathrm{ant}}$

 $\nabla \cdot (\hat{\epsilon} \cdot \nabla \phi) - i k_0 \nabla \cdot (\hat{\epsilon} \cdot \mathbf{A}) = -4\pi \rho_{ant}$

- LEMan: full-wave code / plasma enters with its dielectric tensor using a hot-plasma model accurate to all orders in Larmor radius
- VENUS-LEVIS: particle following in W7-X equilibrium using pre-computed ICRH wave field to apply ICRH kicks to the particles in a Monte-Carlo sense

⁸M. Jucker et al., Comput. Phys. Commun. **183**, 912-925 (2011) IPPIC SLABY ET AL IMAY 20, 2025

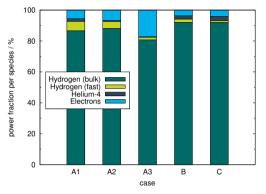
equilibrium

The SCENIC code



What do we use SCENIC for?

- provide theoretical support for ICRH operation at W7-X
- assess power absorption, build-up of a fast-ion distribution function and particle losses
- model "standard heating scenarios" such as the minority heating scenario of H in a ⁴He-plasma (first scenario to be investigated in detail for W7-X)
 - verify that power is indeed absorbed by H-minority
 - check fast-ion generation capabilities of this scheme
- model advanced heating scenarios such as the 3-ion scheme or combined RF-NBI scenarios
- · provide input (profiles, distribution functions, lost-particle data) for other fast-ion codes



Direct RF-power absorption for (H)-⁴He minority scheme in W7-X

F

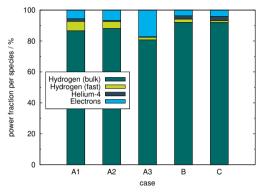
- scan of the minority concentration (5% 15%) performed⁹
- the simulations confirm that most of the power (> 90%) goes to Hydrogen minority
- electrons usually receive the second largest amount of power
- their fraction can increase significantly if too much Hydrogen is in the plasma (case A3)
- very little power goes to the bulk Helium

A*: $n_{\rm e,0} = 5.0 \cdot 10^{19} \ {\rm m}^{-3}$ with increasing minority fraction

B:
$$n_{\rm e,0} = 1.0 \cdot 10^{20} \text{ m}^{-3}$$
 C: $n_{\rm e,0} = 1.5 \cdot 10^{20} \text{ m}^{-3}$

⁹C. Slaby et al., *J. Phys.: Conf. Ser.* **2397**, 012006 (2022) IPPIC. SLABY ET AL. I MAY 20, 2025





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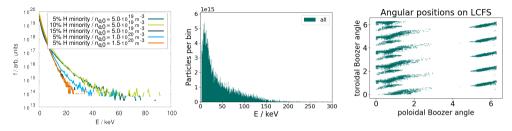
B: $n_{\rm e,0} = 1.0 \cdot 10^{20} \text{ m}^{-3}$ C: $n_{\rm e,0} = 1.5 \cdot 10^{20} \text{ m}^{-3}$

take-away message: regular minority scheme offers good power absorption over a wide range of plasma parameters

⁹C. Slaby et al., *J. Phys.: Conf. Ser.* **2397**, 012006 (2022) IPPIC. SLABY ET AL. I MAY 20, 2025



Fast-ion distribution functions and lost particles

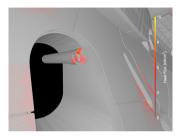


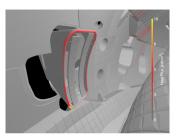
- · H minority absorbs least amount of power at 15% concentration at $n_{\rm e,0} = 5.0 \cdot 10^{19} \text{ m}^{-3} \Longrightarrow$ fewer fast ions (power is also shared among more particles)¹⁰
- SCENIC only simulates the interior of the plasma, $s \leq 1$, but records lost-particle data
- data used to restart particles in an ASCOT simulation which is able to treat the SOL and includes a full model of the 3D wall with all components
- \Rightarrow find hot spots / compare to NBI / assess machine safety aspects

¹⁰C. Slaby et al., *J. Phys.: Conf. Ser.* **2397**, 012006 (2022) IPPIC. SLABY ET AL. I MAY 20, 2025

Wendelstein 7-X

ICRH-induced wall loads





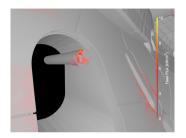
- shown are the wall-loads on the MPM and the ICRH antenna itself
- \cdot power going to the wall is less than for NBI-generated fast ions¹¹ (understandable since ICRH antenna only injects \approx 1 $\rm MW) \rightarrow$ acceptable wall loads¹²
- \cdot simulations show that MPM is hit by fast ions \rightarrow good news for FILDs

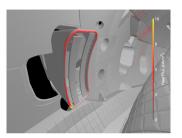
¹²C. Slaby et al., J. Phys.: Conf. Ser. 2397, 012006 (2022)

¹¹J. Kontula et al., *Plasma Phys. Control. Fusion* **65**, 075008 (2023)

Vendelstein 7-X

ICRH-induced wall loads





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- \cdot simulations show that MPM is hit by fast ions \rightarrow good news for FILDs
- · SCENIC-ASCOT coupling established / heat loads to be reassessed for other heating schemes

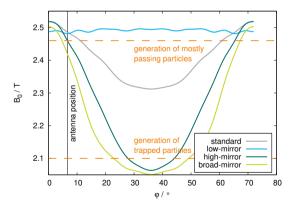
¹¹J. Kontula et al., *Plasma Phys. Control. Fusion* **65**, 075008 (2023) ¹²C. Slaby et al., *J. Phys.: Conf. Ser.* **2397**, 012006 (2022)

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Finding scenarios that favour the generation of deeply trapped fast ions

generating deeply trapped collisionless fast ions (needed to prove the optimisation) is an experimental challenge \Rightarrow NBI does not populate deeply trapped orbits well¹³

- ICRH antenna is located in a region of high field strength \rightarrow mostly passing particles will be generated if resonance is placed in front of the antenna
- ICRH antenna more flexible than NBI \rightarrow change antenna frequency to place the resonance in the triangular plane thus generating deeply trapped fast ions

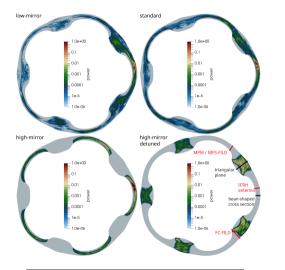


¹³S.A. Lazerson et al., *Phys. Plasmas* **31**, 072506 (2024) IPPIC. SLABY ET AL. I MAY 20, 2025



Shifting the resonance

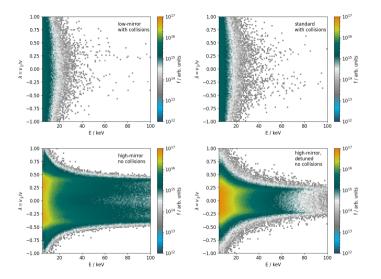




- LEMan simulations confirm that regions of power deposition can indeed be shifted to regions of low *B* (if the mirror ratio of the configuration allows for it)¹⁴
- detuning not possible in low-mirror or standard configuration
- antenna position and locations of some important fast-ion diagnostics indicated in the bottom-right plot
- experimentally: only $f_{\rm ant} = 25 \ {\rm MHz}$ or 37.5 ${\rm MHz}$ were originally foreseen \rightarrow performing such detuned scenarios would require changes to the generators and a re-tuning of the transmission lines
- \Rightarrow a plan for future operation phases

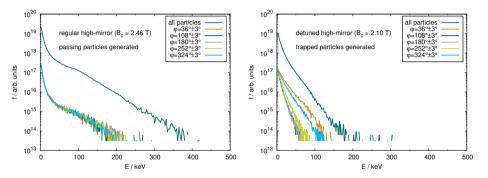
¹⁴C. Slaby et al., accepted by Plasma Phys. Control. Fusion (2025) IPPIC SLABY ET AL. I MAY 20, 2025

Fast-ion distribution functions for detuned scenarios



- collisions rapidly isotropize velocity space and prevent formation of a highly energetic fast-ion distribution function
- without collisions, the beneficial effect of detuning the antenna frequency can be seen more clearly
- ⇒ extent of distribution function in λ -direction becomes more narrow
 - simulations with collisions but with higher ${\cal T}_{\rm e}$ (increases slowing-down time) \Rightarrow they fall between the extremes shown here

Toroidal asymmetry of the fast-ion distribution function (collisionless)



· if passing particles are generated, the distribution function is identical at every toroidal angle

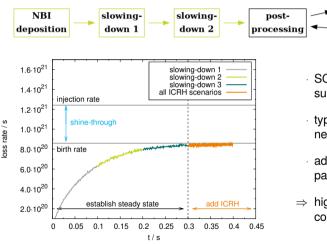
· if predominantly trapped fast ions are generated, the particles cannot leave the field period in which they are generated in and more fast ions reside in modules close the ICRH antenna \rightarrow detectors should be placed there (much better signal-to-noise ratio)¹⁵

¹⁵C. Slaby et al., *accepted by Plasma Phys. Control. Fusion* (2025) IPPIC. SLABY ET AL. I MAY 20, 2025



combined RF-NBI

Combined RF-NBI scenarios



SCENIC can model NBI deposition and the subsequent slowing-down of the NBI ions^{16,17}

merge

wave fields

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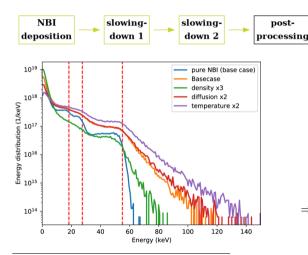
- typically several slowing-down simulations necessary until steady-state is reached
- adding ICRH accelerates the already fast particles further
- \Rightarrow higher energies (> 100 keV) can be reached compared to using only NBI

¹⁶M. Machielsen et al., *J. Plasma Phys.* 89, 955890202 (2023)
 ¹⁷C. Slaby et al., *20th European Fusion Theory Conference (EFTC)*, Padua, Italy (2023)
 ¹⁷I. Slaby et al., *1May 20, 2025*



combined RF-NBI

Combined RF-NBI scenarios



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wave fields

LEMan

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Summary and conclusions



- · ICRH physics at W7-X is investigated with the SCENIC code
 - direct RF-power absorption, fast-ion generation, trapped and passing, RF-NBI
- also simpler tools than SCENIC are in use for designing and analyzing experiments with ICRH
 ICRH overview tool
- see poster Wednesday-23 by D. Hartmann for first experimental results with the ICRH antenna at W7-X

Ongoing and future work

- development of a new ICRH full wave code based on integral kernels and finite elements instead of Fourier expansions (ENR project led by LPP ERM-KMS Brussels)
 - better coupling to antenna models (e.g. PETRA-M), local mesh refinement
 - see also I-16 by P.U. Lamalle and poster Tuesday-21 by B. Reman for the theory behind the approach and first numerical results